

TOP-BESS MODEL AND ITS PHENOMENOLOGY

M. Gintner^{1,2}, J. Juráň², I. Melo¹

¹ Žilina U.

² IEAP CTU Prague

MFF UK Bratislava, Nov 15, 2011

putting things straight beforehand ...

top-BESS model



NO

- fundamental theory
- ESB mechanism



YES

- effective Lagrangian
- strong ESB physics

putting things straight beforehand ...

top-BESS model



NO

- fundamental theory
- ESB mechanism



YES

- effective Lagrangian
- strong ESB physics

putting things straight beforehand ...

top-BESS model



NO

- fundamental theory
- ESB mechanism



YES

- effective Lagrangian
- strong ESB physics

OUTLINE

- 1 INTRODUCTION
- 2 TOP-BESS MODEL
- 3 PHENOMENOLOGY

OUTLINE

- 1 INTRODUCTION
- 2 TOP-BESS MODEL
- 3 PHENOMENOLOGY

GAUGE PRINCIPLE

⇒ Gauge Bosons

⇒ Interactions

⇒ Renormalizability

GAUGE PRINCIPLE

⇒ Gauge Bosons

⇒ Interactions

⇒ Renormalizability

GAUGE PRINCIPLE

⇒ Gauge Bosons

⇒ Interactions

⇒ Renormalizability

GAUGE PRINCIPLE

⇒ Gauge Bosons

⇒ Interactions

⇒ Renormalizability

SM \rightarrow GP SUCCESS

$$\text{SM: } SU(3)_C \times SU(2)_L \times U(1)_Y$$

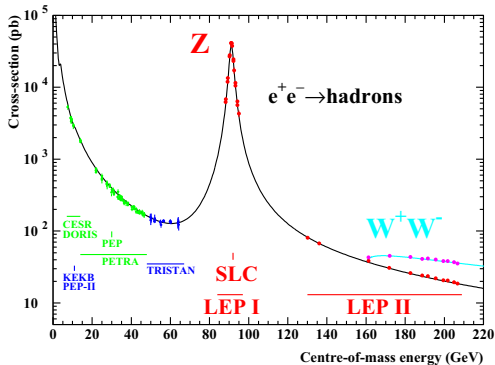
gluons

W^\pm, Z

photon

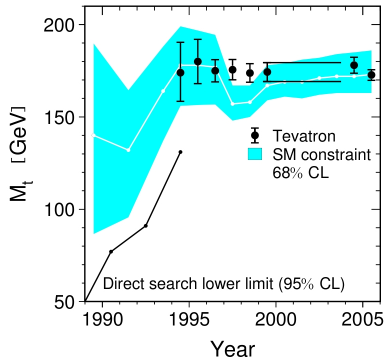
SM \rightarrow GP SUCCESS

SM: $SU(3)_C \times SU(2)_L \times U(1)_Y$



SM \rightarrow GP SUCCESS

SM: $SU(3)_C \times SU(2)_L \times U(1)_Y$



EW SYMMETRY BROKEN!

- fermions:

$$m_t \gg m_b > \dots > m_e \gg m_{\nu_e}$$

- weak gauge bosons:

$$M_Z, M_W \approx 100 \text{ GeV}$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int}$$

EW SYMMETRY BROKEN!

- fermions:

$$m_t \gg m_b > \dots > m_e \gg m_{\nu_e}$$

- weak gauge bosons:

$$M_Z, M_W \approx 100 \text{ GeV}$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int}$$

EW SYMMETRY BROKEN!

- fermions:

$$m_t \gg m_b > \dots > m_e \gg m_{\nu_e}$$

- weak gauge bosons:

$$M_Z, M_W \approx 100 \text{ GeV}$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int} + \frac{1}{2} M_{GB}^2 X_\mu X^\mu + m_f (\bar{\psi}_L \psi_R + \text{h.c.})$$

EW SYMMETRY BROKEN!

- fermions:

$$m_t \gg m_b > \dots > m_e \gg m_{\nu_e}$$

- weak gauge bosons:

$$M_Z, M_W \approx 100 \text{ GeV}$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int} + \underbrace{\frac{1}{2}M_{GB}^2 X_\mu X^\mu + m_f(\bar{\psi}_L \psi_R + \text{h.c.})}_{SU(2)_L \times U(1)_Y \text{ broken!}}$$

SAVING THE GAUGE PRINCIPLE

ESB = **spontaneous** symmetry breaking:

$$\text{symm}(\textit{vacuum}) < \text{symm}(\textit{Lagr})$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int} + \mathcal{L}_{SSB}$$

$$SU(2)_L \times U(1)_Y \xrightarrow{SSB} U(1)_{em}$$

3 Goldstone bosons

$$\mathcal{L}_{SSB} = ?$$

SAVING THE GAUGE PRINCIPLE

ESB = **spontaneous** symmetry breaking:

$$\text{symm}(\textit{vacuum}) < \text{symm}(\textit{Lagr})$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int} + \mathcal{L}_{SSB}$$

$$SU(2)_L \times U(1)_Y \xrightarrow{SSB} U(1)_{em}$$

3 Goldstone bosons

$$\mathcal{L}_{SSB} = ?$$

SAVING THE GAUGE PRINCIPLE

ESB = **spontaneous** symmetry breaking:

$$\text{symm}(\textit{vacuum}) < \text{symm}(\textit{Lagr})$$

$$\mathcal{L} = \mathcal{L}_{kin} + \mathcal{L}_{int} + \mathcal{L}_{SSB}$$

$$SU(2)_L \times U(1)_Y \xrightarrow{SSB} U(1)_{em}$$

3 Goldstone bosons

$$\mathcal{L}_{SSB} = ?$$

BENCHMARK HYPOTHESIS \rightarrow SM HIGGS

- $SU(2)_L$ complex scalar doublet Φ

- $v = \langle 0|\Phi|0\rangle$

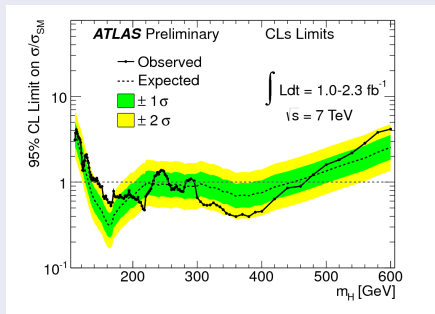
Higgs boson

BENCHMARK HYPOTHESIS → SM HIGGS

- $SU(2)_L$ complex scalar doublet Φ

- $v = \langle 0|\Phi|0\rangle$

Higgs boson

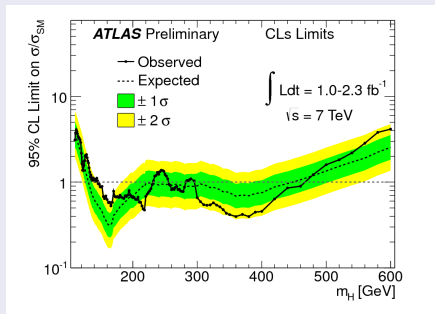


BENCHMARK HYPOTHESIS → SM HIGGS

- $SU(2)_L$ complex scalar doublet Φ

- $v = \langle 0|\Phi|0\rangle$

Higgs boson



theoretical pathologies



just a parameterization
of ESB?

HIGGS BOSON ALTERNATIVES

heavy/no Higgs **violates unitarity** ≈ 1 TeV

weakly interacting:

- new forces and particles
- *perturbative*
- more Higgses, SUSY

strongly interacting:

- new forces and particles
- *non-perturbative* \rightarrow *bound states*
- TC and its extensions



extra-dimensions:

4D *strongly* interacting



5D *weakly* interacting

HIGGS BOSON ALTERNATIVES

heavy/no Higgs **violates unitarity** ≈ 1 TeV

weakly interacting:

- new forces and particles
- *perturbative*
- more Higgses, SUSY

strongly interacting:

- new forces and particles
- *non-perturbative* \rightarrow *bound states*
- TC and its extensions

extra-dimensions:

4D *strongly* interacting



5D *weakly* interacting

HIGGS BOSON ALTERNATIVES

heavy/no Higgs **violates unitarity** ≈ 1 TeV

weakly interacting:

- new forces and particles
- *perturbative*
- more Higgses, SUSY

strongly interacting:

- new forces and particles
- *non-perturbative* \rightarrow *bound states*
- TC and its extensions

extra-dimensions:

4D *strongly* interacting



5D *weakly* interacting

HIGGS BOSON ALTERNATIVES

heavy/no Higgs **violates unitarity** ≈ 1 TeV

weakly interacting:

- new forces and particles
- *perturbative*
- more Higgses, SUSY



strongly interacting:

- new forces and particles
- *non-perturbative* \rightarrow *bound states*
- TC and its extensions



extra-dimensions:

4D *strongly* interacting



5D *weakly* interacting

OUTLINE

- 1 INTRODUCTION
- 2 TOP-BESS MODEL
- 3 PHENOMENOLOGY

EFFECTIVE DESCRIPTION OF STRONG ESB

$SU(2)_L \times U(1)_Y$ broken *dynamically*:

- *not* solvable perturbatively
- *chiral effective* Lagrangian for **Goldstone bosons**

nonlinear sigma model

$$\mathcal{L} = \frac{v^2}{2} \text{Tr} \left[(\partial_\mu U^\dagger) (\partial^\mu U) \right]$$

$$U = \exp(2i\pi^a \tau^a / v)$$

- ... + **resonances**

scalar, vector, ...

LHC \rightarrow the *lightest* BSM resonances

EFFECTIVE DESCRIPTION OF STRONG ESB

$SU(2)_L \times U(1)_Y$ broken *dynamically*:

- *not* solvable perturbatively
- *chiral effective* Lagrangian for Goldstone bosons

nonlinear sigma model

$$\mathcal{L} = \frac{v^2}{2} \text{Tr} [(\partial_\mu U^\dagger)(\partial^\mu U)]$$

$$U = \exp(2i\pi^a \tau^a / v)$$

- ... + resonances

scalar, vector, ...

LHC \rightarrow the *lightest* BSM resonances

EFFECTIVE DESCRIPTION OF STRONG ESB

$SU(2)_L \times U(1)_Y$ broken *dynamically*:

- *not* solvable perturbatively
- *chiral effective* Lagrangian for **Goldstone bosons**

nonlinear sigma model

$$\mathcal{L} = \frac{v^2}{2} \text{Tr} \left[(\partial_\mu U^\dagger)(\partial^\mu U) \right]$$

$$U = \exp(2i\pi^a \tau^a / v)$$

- ... + **resonances**

scalar, vector, ...

LHC \rightarrow the *lightest* BSM resonances

EFFECTIVE DESCRIPTION OF STRONG ESB

$SU(2)_L \times U(1)_Y$ broken *dynamically*:

- *not* solvable perturbatively
- *chiral effective* Lagrangian for **Goldstone bosons**

nonlinear sigma model

- ... + **resonances**

scalar, vector, ...

$$\mathcal{L} = \frac{v^2}{2} \text{Tr} \left[(\partial_\mu U^\dagger)(\partial^\mu U) \right]$$

$$U = \exp(2i\pi^a \tau^a / v)$$

LHC \rightarrow the *lightest* BSM resonances

EFFECTIVE DESCRIPTION OF STRONG ESB

$SU(2)_L \times U(1)_Y$ broken *dynamically*:

- *not* solvable perturbatively
- *chiral effective* Lagrangian for **Goldstone bosons**

nonlinear sigma model

$$\mathcal{L} = \frac{v^2}{2} \text{Tr} \left[(\partial_\mu U^\dagger)(\partial^\mu U) \right]$$

$$U = \exp(2i\pi^a \tau^a / v)$$

- ... + **resonances**

scalar, vector, ...

LHC \rightarrow the *lightest* BSM resonances

HIDDEN LOCAL SYMMETRY

M. Bando, T. Kugo, K. Yamawaki (1984)

Any $NL\sigma M(G/H)$ is gauge equivalent to "linear" $G_{glob} \times H_{loc}$ model. $G_{glob} \times H_{loc}$ LAGRANGIAN... $g \in G, h(x) \in H_{loc}$

$$\xi = e^{i\pi^a X^a/v} e^{i\sigma^b S^b/v} \longrightarrow g \xi h(x)$$

$$V_\mu = V_\mu^a S^a \longrightarrow h^\dagger V_\mu h + h^\dagger \partial_\mu h$$

$$\omega_\mu = \xi^\dagger \partial_\mu \xi \in G, \quad \omega_\mu^{\parallel,\perp} = [\omega_\mu \pm \tau(\omega_\mu)]/2 \in H, G \setminus H$$

$$\mathcal{L} = -v^2 \left[\text{Tr}(\omega_\mu^\perp)^2 + \alpha \text{Tr}(\omega_\mu^\parallel - V_\mu)^2 \right] = \frac{v^2}{4} \text{Tr}[(\partial_\mu U^\dagger)(\partial^\mu U)]$$

$$\omega_\mu^\perp = \xi^\dagger (\partial_\mu U) \tau(\xi)/2 \quad \dots \quad U = \xi \tau(\xi^\dagger)$$

HIDDEN LOCAL SYMMETRY

M. Bando, T. Kugo, K. Yamawaki (1984)

Any $NL\sigma M(G/H)$ is gauge equivalent to "linear" $G_{glob} \times H_{loc}$ model. $G_{glob} \times H_{loc}$ LAGRANGIAN... $g \in G, h(x) \in H_{loc}$

$$\xi = e^{i\pi^a X^a/v} e^{i\sigma^b S^b/v} \longrightarrow g \xi h(x)$$

$$\mathbf{V}_\mu = V_\mu^a S^a \longrightarrow h^\dagger \mathbf{V}_\mu h + h^\dagger \partial_\mu h$$

$$\omega_\mu = \xi^\dagger \partial_\mu \xi \in \mathcal{G}, \quad \omega_\mu^{\parallel, \perp} = [\omega_\mu \pm \tau(\omega_\mu)]/2 \in \mathcal{H}, \mathcal{G} \setminus \mathcal{H}$$

$$\mathcal{L} = -v^2 \left[\text{Tr}(\omega_\mu^\perp)^2 + \alpha \text{Tr}(\omega_\mu^\parallel - \mathbf{V}_\mu)^2 \right] = \frac{v^2}{4} \text{Tr}[(\partial_\mu \mathbf{U}^\dagger)(\partial^\mu \mathbf{U})]$$

$$\omega_\mu^\perp = \xi^\dagger (\partial_\mu \mathbf{U}) \tau(\xi)/2 \quad \dots \mathbf{U} = \xi \tau(\xi^\dagger)$$

TRANSITION TO “U-GAUGE”

$$\nearrow h(x) = e^{-i\sigma(x)/v} \in H_{loc} \searrow$$

linear

$$\begin{aligned} \xi(x) &= e^{i\pi(x)/v} e^{i\sigma(x)/v} \\ U = \xi \tau(\xi^\dagger) &= e^{2i\pi(x)/v} \end{aligned}$$

$$G \times H_{loc} : \quad \dots h(x) \in H_{loc}$$

$$\begin{aligned} \xi &\longrightarrow g \xi h(x) \\ U &\longrightarrow g U \tau(g^\dagger) \\ \mathbf{V}_\mu &\longrightarrow h^\dagger(x) \mathbf{V}_\mu h(x) + h^\dagger \partial_\mu h \end{aligned}$$

$$\mathcal{L} = \frac{v^2}{4} \text{Tr}[(\partial_\mu U^\dagger)(\partial^\mu U)]$$

U-gauge:

non-linear

$$\begin{aligned} \xi(x) &= e^{i\pi(x)/v} \\ U = \xi \tau(\xi^\dagger) &= e^{2i\pi(x)/v} \end{aligned}$$

$$G : \quad \dots g_h \in H \subset G$$

$$\begin{aligned} \xi &\longrightarrow g \xi g_h^\dagger(g, \xi) \\ U &\longrightarrow g U \tau(g^\dagger) \\ \mathbf{V}_\mu &\longrightarrow g_h(g, \xi) \mathbf{V}_\mu g_h^\dagger(g, \xi) + g_h \partial_\mu g_h^\dagger \end{aligned}$$

$$\mathcal{L} = \frac{v^2}{4} \text{Tr}[(\partial_\mu U^\dagger)(\partial^\mu U)]$$

BESS MODEL

BREAKING ELECTROWEAK SYMMETRY STRONGLY

R. Casalbuoni, S. De Curtis, D. Dominici, R. Gatto
PLB155, 95 (1985), NPB282, 235 (1987)

- effective Lagrangian
- \mathcal{H}_{SM} + new vector resonances

$$\mathcal{L}_{BESS} = \mathcal{L}_{GB}(W, B, V) + \mathcal{L}_{ESB}(\vec{\pi}, \vec{\sigma}) + \mathcal{L}_{ferm}$$

BESS MODEL

BREAKING ELECTROWEAK SYMMETRY STRONGLY

R. Casalbuoni, S. De Curtis, D. Dominici, R. Gatto
PLB**155**, 95 (1985), NPB**282**, 235 (1987)

- effective Lagrangian
- $\mathbb{H}SM$ + new vector resonances

$$\mathcal{L}_{BESS} = \mathcal{L}_{GB}(W, B, V) + \mathcal{L}_{ESB}(\vec{\pi}, \vec{\sigma}) + \mathcal{L}_{ferm}$$

BESS SYMMETRIES AND COUPLINGS

- *global symmetry:*

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times SU(2)_{HLS} \xrightarrow{SSB} SU(2)_{L+R} \times U(1)_{B-L}$$

- *local symmetry:*

$$\frac{SU(2)_L \times U(1)_Y \times SU(2)_{HLS}}{g \quad g' \quad g''} \xrightarrow{SSB} \frac{U(1)_{em}}{e}$$

- *gauge sector:*

$$W^\pm, Z \quad A \quad V^\pm, V^0 \quad \dots \text{mixing}$$

- *fermion sector:*

◇ direct cplg: ... $bg'' \bar{\psi}_L \mathcal{Y} \psi_L, b'g'' \bar{\psi}_R \mathcal{Y} \psi_R$... **universal**

◇ indirect cplg: ... $1/g'' \bar{\psi}(\mathcal{Z}, W)\psi$... mixing induced

BESS SYMMETRIES AND COUPLINGS

- *global symmetry:*

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times SU(2)_{HLS} \xrightarrow{SSB} SU(2)_{L+R} \times U(1)_{B-L}$$

- *local symmetry:*

$$\begin{array}{ccc} SU(2)_L \times U(1)_Y \times SU(2)_{HLS} & \xrightarrow{SSB} & U(1)_{em} \\ g & g' & g'' & e \end{array}$$

- *gauge sector:*

$$W^\pm, Z \quad A \quad V^\pm, V^0 \quad \dots \text{mixing}$$

- *fermion sector:*

◇ direct cplg: ... $bg'' \bar{\psi}_L \not{Y} \psi_L, b'g'' \bar{\psi}_R \not{Y} \psi_R$... **universal**

◇ indirect cplg: ... $1/g'' \bar{\psi}(\not{Z}, \not{W})\psi$... mixing induced

BESS SYMMETRIES AND COUPLINGS

- *global symmetry:*

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times SU(2)_{HLS} \xrightarrow{SSB} SU(2)_{L+R} \times U(1)_{B-L}$$

- *local symmetry:*

$$\begin{array}{ccc} SU(2)_L \times U(1)_Y \times SU(2)_{HLS} & \xrightarrow{SSB} & U(1)_{em} \\ g \quad \quad \quad g' \quad \quad \quad g'' & & e \end{array}$$

- *gauge sector:*

$$W^\pm, Z \quad A \quad V^\pm, V^0 \quad \dots \text{mixing}$$

- *fermion sector:*

◇ direct cplg: ... $bg'' \bar{\psi}_L \mathcal{Y} \psi_L, b'g'' \bar{\psi}_R \mathcal{Y} \psi_R$... universal

◇ indirect cplg: ... $1/g'' \bar{\psi}(\mathcal{Z}, W)\psi$... mixing induced

BESS SYMMETRIES AND COUPLINGS

- *global symmetry:*

$$SU(2)_L \times SU(2)_R \times U(1)_{B-L} \times SU(2)_{HLS} \xrightarrow{SSB} SU(2)_{L+R} \times U(1)_{B-L}$$

- *local symmetry:*

$$\begin{array}{ccc} SU(2)_L \times U(1)_Y \times SU(2)_{HLS} & \xrightarrow{SSB} & U(1)_{em} \\ g \quad g' \quad g'' & & e \end{array}$$

- *gauge sector:*

$$W^\pm, Z \quad A \quad V^\pm, V^0 \quad \dots \text{mixing}$$

- *fermion sector:*

◇ direct cplg: ... $bg'' \bar{\psi}_L \not{Y} \psi_L, b'g'' \bar{\psi}_R \not{Y} \psi_R$... **universal**

◇ indirect cplg: ... $1/g'' \bar{\psi}(\not{Z}, \not{W})\psi$... mixing induced

OUTSTANDING TOP QUARK

$$m_t \approx v/\sqrt{2} \quad \rightarrow \quad \text{special role in ESB?}$$

new physics behind m_t

ESB related

Extended TC, ...

ESB *unrelated*

Topcolor Assisted TC, ...

OUTSTANDING TOP QUARK

$m_t \approx v/\sqrt{2}$ → special role in ESB?

new physics behind m_t

↙
ESB related

Extended TC, ...

↘
ESB *unrelated*

Topcolor Assisted TC, ...

- *gauge sector* \equiv *BESS*

- *fermion sector*:

- 3rd quark generation only

... b_L, b_R

- $bottom_R$ vs. top_R

... p

- new fermion terms

... λ_L, λ_R

- *gauge sector* \equiv *BESS*

- *fermion sector*:

- ◊ 3rd quark generation only

... b_L, b_R

- ◊ $bottom_R$ vs. top_R

... p

- ◊ new fermion terms

... λ_L, λ_R

- *gauge sector* \equiv *BESS*

- *fermion sector*:

- ◊ 3rd quark generation only

... b_L, b_R

- ◊ $bottom_R$ vs. top_R

... p

- ◊ new fermion terms

... λ_L, λ_R

- *gauge sector* \equiv *BESS*

- *fermion sector*:

- ◊ 3rd quark generation only

... b_L, b_R

- ◊ *bottom*_R vs. *top*_R

... p

- ◊ new fermion terms

... λ_L, λ_R

- *gauge sector* \equiv BESS
- *fermion sector*:
 - ◇ 3rd quark generation only ... b_L, b_R
 - ◇ $bottom_R$ vs. top_R ... p
 - ◇ new fermion terms ... λ_L, λ_R

- *gauge sector* \equiv BESS
- *fermion sector*:
 - ◇ 3rd quark generation only ... b_L, b_R
 - ◇ $bottom_R$ vs. top_R ... p
 - ◇ new fermion terms ... λ_L, λ_R

OUTLINE

- 1 INTRODUCTION
- 2 TOP-BESS MODEL
- 3 PHENOMENOLOGY

NEW RESONANCE MASSES

- mass of the vector resonance:

$$M_V = \frac{\sqrt{\alpha} g'' v}{2}$$

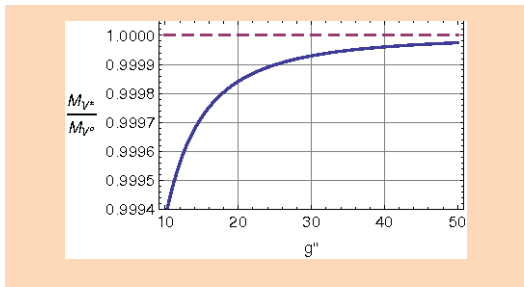
- EW gauge bosons \rightarrow mixing \rightarrow mass splitting

NEW RESONANCE MASSES

- mass of the vector resonance:

$$M_V = \frac{\sqrt{\alpha} g'' v}{2}$$

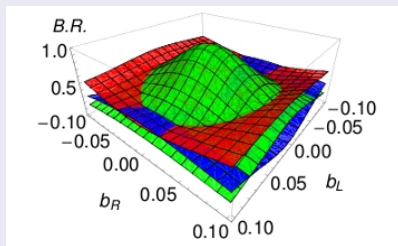
- EW gauge bosons \rightarrow mixing \rightarrow mass splitting



DECAY WIDTHS

- $V^0 \rightarrow W^+W^- + t\bar{t} + b\bar{b} + \dots$
- $V^+ \rightarrow W^+Z + t\bar{b} + \dots$
- $\Gamma \sim 10 - 100 \text{ GeV}$

$$M_V = 1 \text{ TeV}, g'' = 20, p = 0$$



$WW, t\bar{t}, b\bar{b}$

UNITARITY CONSTRAINTS

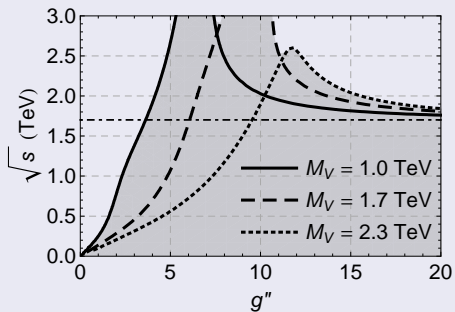
- GB scatterings:

$$W_L^+ W_L^-, Z_L Z_L,$$

$$W_L^\pm Z_L, W_L^\pm W_L^\pm$$

- tree level

- Equivalence Theorem



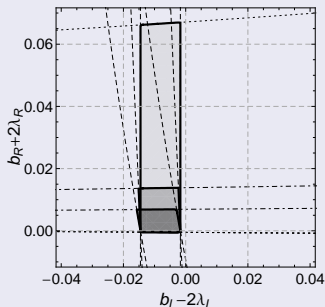
LOW-ENERGY LIMITS

EXPERIMENT: LEP + SLC + TEVATRON

EWPD ϵ -analysis: $\epsilon_1, \epsilon_2, \epsilon_3, \epsilon_b, \Gamma(Z \rightarrow b\bar{b}), B \rightarrow X_s\gamma, p\bar{p} \rightarrow WZX$

LOW-ENERGY LIMITS

EXPERIMENT: LEP + SLC + TEVATRON

EWPD ϵ -analysis: $\epsilon_1, \epsilon_2, \epsilon_3, \epsilon_b, \Gamma(Z \rightarrow b\bar{b}), B \rightarrow X_s\gamma, p\bar{p} \rightarrow WZX$ Intersections of 90% C.L.
allowed regions.

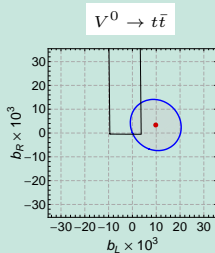
$$M_V = 1 \text{ TeV}$$
$$g'' = 10$$

THE DEATH VALLEY

direct + indirect *cplgs* \Rightarrow *DV*

THE DEATH VALLEY

direct + indirect *cplgs* \Rightarrow *DV*

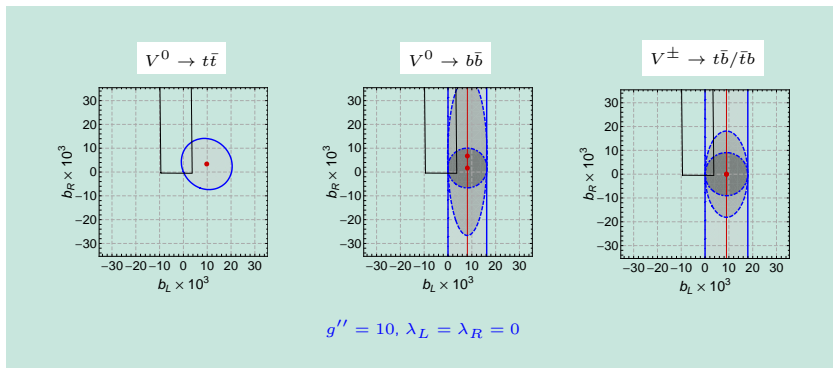


$$g'' = 10, \lambda_L = \lambda_R = 0$$

The Death Valley regions of the $V \rightarrow t\bar{t}/b\bar{b}/tb$ decays.

THE DEATH VALLEY

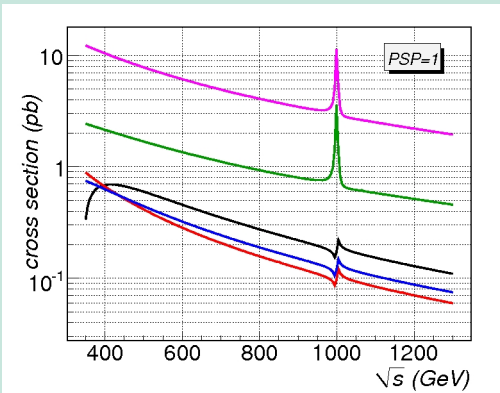
direct + indirect *cplgs* \Rightarrow *DV*



The Death Valley regions of the $V \rightarrow t\bar{t}/b\bar{b}/tb$ decays.

HIDING THE PEAK

$$M_V = 1 \text{ TeV}, g'' = 20, p = 0, \lambda_R = 0$$

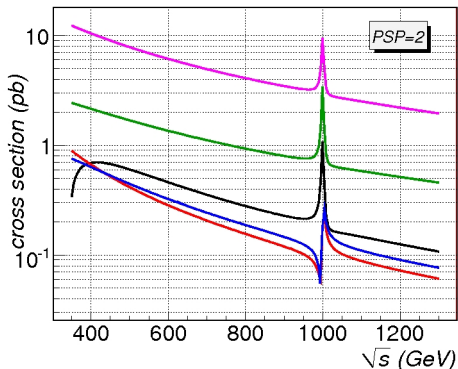


- **no direct cplng**
 $b_L = 0$
 $b_R = 0$
 $\lambda_L = 0$
- **outside the DV**
 $b_L = -0.010$
 $b_R = +0.030$
 $\lambda_L = 0$
- **$t\bar{b}$ & $b\bar{b}$ in the DV**
 $b_L = +0.009$
 $b_R = +0.030$
 $\lambda_L = +0.006$
- **all in the DV**
 $b_L = +0.0098$
 $b_R = +0.0034$
 $\lambda_L = +0.006$

$e^+e^- \rightarrow W^+W^-$
 $u\bar{d} \rightarrow W^+Z$
 $e^+e^- \rightarrow t\bar{t}$
 $u\bar{d} \rightarrow t\bar{b}$
 $e^+e^- \rightarrow b\bar{b}$

HIDING THE PEAK

$$M_V = 1 \text{ TeV}, g'' = 20, p = 0, \lambda_R = 0$$



- no direct cplng
 $b_L = 0$
 $b_R = 0$
 $\lambda_L = 0$
- outside the DV
 $b_L = -0.010$
 $b_R = +0.030$
 $\lambda_L = 0$
- $t\bar{b}$ & $b\bar{b}$ in the DV
 $b_L = +0.009$
 $b_R = +0.030$
 $\lambda_L = +0.006$
- all in the DV
 $b_L = +0.0098$
 $b_R = +0.0034$
 $\lambda_L = +0.006$

$$e^+e^- \rightarrow W^+W^-$$

$$u\bar{d} \rightarrow W^+Z$$

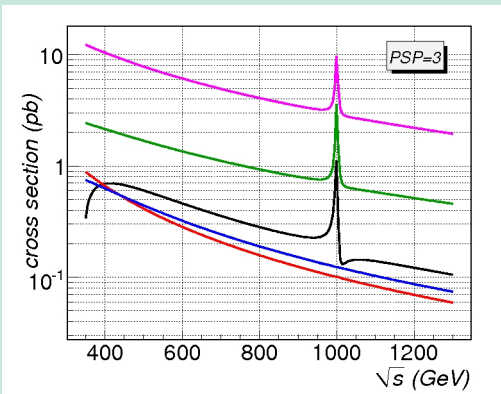
$$e^+e^- \rightarrow t\bar{t}$$

$$u\bar{d} \rightarrow t\bar{b}$$

$$e^+e^- \rightarrow b\bar{b}$$

HIDING THE PEAK

$$M_V = 1 \text{ TeV}, g'' = 20, p = 0, \lambda_R = 0$$

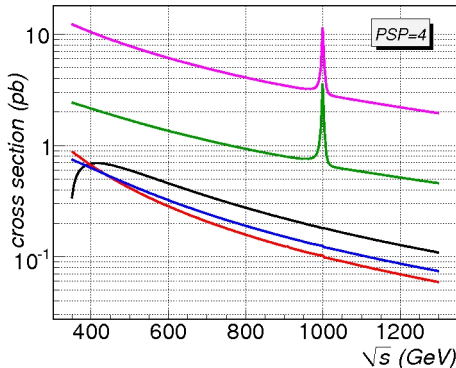


- no direct cplng
 $b_L = 0$
 $b_R = 0$
 $\lambda_L = 0$
- outside the DV
 $b_L = -0.010$
 $b_R = +0.030$
 $\lambda_L = 0$
- $t\bar{b}$ & $b\bar{b}$ in the DV
 $b_L = +0.009$
 $b_R = +0.030$
 $\lambda_L = +0.006$
- all in the DV
 $b_L = +0.0098$
 $b_R = +0.0034$
 $\lambda_L = +0.006$

$e^+e^- \rightarrow W^+W^-$
 $u\bar{d} \rightarrow W^+Z$
 $e^+e^- \rightarrow t\bar{t}$
 $u\bar{d} \rightarrow t\bar{b}$
 $e^+e^- \rightarrow b\bar{b}$

HIDING THE PEAK

$$M_V = 1 \text{ TeV}, g'' = 20, p = 0, \lambda_R = 0$$



- no direct cplng
 $b_L = 0$
 $b_R = 0$
 $\lambda_L = 0$
- outside the DV
 $b_L = -0.010$
 $b_R = +0.030$
 $\lambda_L = 0$
- $t\bar{b}$ & $b\bar{b}$ in the DV
 $b_L = +0.009$
 $b_R = +0.030$
 $\lambda_L = +0.006$
- all in the DV
 $b_L = +0.0098$
 $b_R = +0.0034$
 $\lambda_L = +0.006$

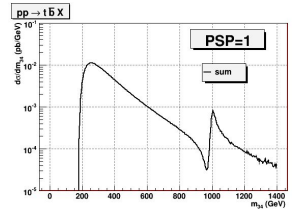
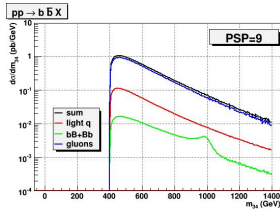
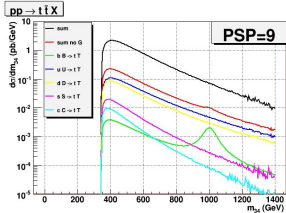
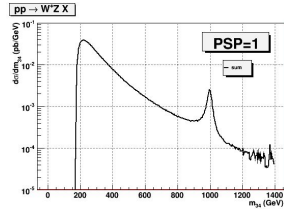
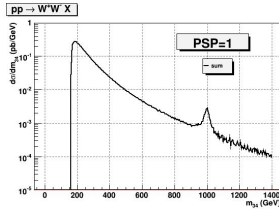
$e^+e^- \rightarrow W^+W^-$
 $u\bar{d} \rightarrow W^+Z$
 $e^+e^- \rightarrow t\bar{t}$
 $u\bar{d} \rightarrow t\bar{b}$
 $e^+e^- \rightarrow b\bar{b}$

WHAT'S NEXT?

- theoretical development
 - low-energy limits
 - scrutinizing the parameter space
 - relation to existing theories
 - ...
- probing tBESS at LHC and ILC
 - Drell-Yan processes at LHC
 - ...

DRELL-YAN AT LHC

... PEEKING



CONCLUSIONS

- effective description of strong ESB new physics needed
- top-BESS — modification of BESS, special role of top quark
 - new $SU(2)$ resonance triplet
 - direct coupling to top and bottom
 - λ -terms
- low-E limits on the fermion parameters relaxed
- the Death Valley effect
- LHC: Drell-Yan processes

CONCLUSIONS

- effective description of strong ESB new physics needed
- top-BESS — modification of BESS, special role of top quark
 - ◇ *new $SU(2)$ resonance triplet*
 - ◇ *direct coupling to top and bottom*
 - ◇ *λ -terms*
- low-E limits on the fermion parameters relaxed
- the Death Valley effect
- LHC: Drell-Yan processes

CONCLUSIONS

- effective description of strong ESB new physics needed
- top-BESS — modification of BESS, special role of top quark
 - ◇ *new $SU(2)$ resonance triplet*
 - ◇ *direct coupling to top and bottom*
 - ◇ *λ -terms*
- low-E limits on the fermion parameters relaxed
- the Death Valley effect
- LHC: Drell-Yan processes

CONCLUSIONS

- effective description of strong ESB new physics needed
- top-BESS — modification of BESS, special role of top quark
 - ◇ *new $SU(2)$ resonance triplet*
 - ◇ *direct coupling to top and bottom*
 - ◇ *λ -terms*
- low-E limits on the fermion parameters relaxed
- the Death Valley effect
- LHC: Drell-Yan processes

CONCLUSIONS

- effective description of strong ESB new physics needed
- top-BESS — modification of BESS, special role of top quark
 - ◇ *new $SU(2)$ resonance triplet*
 - ◇ *direct coupling to top and bottom*
 - ◇ *λ -terms*
- low-E limits on the fermion parameters relaxed
- the Death Valley effect
- LHC: Drell-Yan processes

I WANT YOU !!!



Enlist Now!